## CIISC5706 Topics in Thenertical Computer Science



Instructor: Shengyu Zhang

## Roadmap

- Intro to math model of quantum mechanics
- Review of quantum algorithms
- The power of quantum computers.
- Quantum games.


## Postulate 1: States

- State space: Every isolated physical system corresponds to a unit vector in a complex vector space.
- Unit vector: $\ell_{2}$-norm is 1.
- Such states are called pure states.
- We use a weird "ket" notation $|\cdot\rangle$ to denote such a state.


## Ket notation

- Mathematically,|•> is a column vector.
- And $\langle\cdot|$ is a row vector.
- $\langle\psi \mid \phi\rangle$ is the inner product between the vectors $|\phi\rangle$ and $|\psi\rangle$.
- $\langle\psi| M|\psi\rangle$ is just the quadratic form $\psi^{T} M \psi$.
- A quantum bit, or qubit, is a state of the form

$$
\alpha|0\rangle+\beta|1\rangle
$$

where $\alpha, \beta \in \mathbb{C}$ are called amplitudes, satisfying that $|\alpha|^{2}+|\beta|^{2}=1$.

- So a qubit can sit anywhere between 0 and 1 (on the unit circle).
- We say that the state is in superposition of $|0\rangle$ and $|1\rangle$.


A quantum bit (qubit)

## Postulate 2: operation

- Evolution: The evolution of a closed quantum system is described by a unitary transformation.
- That is, if a system is in state $\left|\psi_{1}\right\rangle$ at time $t_{1}$, and in state $\left|\psi_{2}\right\rangle$ at time $t_{2}$, then there is a unitary transformation $U$ s.t.

$$
\left|\psi_{2}\right\rangle=U\left|\psi_{1}\right\rangle .
$$

- Unitary transformation: $U^{\dagger}=U^{-1}$
- Recall: $U^{\dagger}=\left(U^{T}\right)^{*}$, transpose + complex conjugate
- You can think of it as a rotation operation.


## Postulate 3: measurement

- Measurement: We can only observe a quantum system by measuring it.
- The outcome of the measurement is random.
- And the system is changed by the measurement.
- If we measure qubit $\alpha|0\rangle+\beta|1\rangle$ in the computational basis $\{|0\rangle,|1\rangle\}$, then outcome " 0 " occurs with prob. $|\alpha|^{2}$, and outcome "1" occurs with prob. $|\beta|^{2}$.
- The system becomes $|0\rangle$ if outcome " 0 " occurs, and becomes $|1\rangle$ if outcome " 1 "


A quantum bit (qubit) occurs.

- The system collapses.


## Measurement on general states

- In general, an orthogonal measurement of a $d$-dim state is given by an orthonormal basis $\left\{\left|\psi_{1}\right\rangle, \ldots\left|\psi_{d}\right\rangle\right\}$.
- If we measure state $|\phi\rangle$ in basis
$\left\{\left|\psi_{1}\right\rangle, \ldots\left|\psi_{d}\right\rangle\right\}$, then outcome $i \in\{1, \ldots, d\}$ occurs with prob. $\left|\left\langle\phi \mid \psi_{i}\right\rangle\right|^{2}$.
- The system collapses to $\left|\psi_{i}\right\rangle$ if outcome $i$ occurs.


## Postulate 4: composition

- Composition: The state of the joint system $\left(S_{1}, S_{2}\right)$, where $S_{1}$ is in state $\left|\psi_{1}\right\rangle$ and $S_{2}$ in $\left|\psi_{2}\right\rangle$, is $\left|\psi_{1}\right\rangle \otimes\left|\psi_{2}\right\rangle$.
- $\otimes$ : tensor product of vectors.
$-\left(a_{1}, a_{2}\right) \otimes\left(b_{1}, b_{2}, b_{3}\right)=\left(a_{1} b_{1}, a_{1} b_{2}, a_{1} b_{3}, a_{2} b_{1}, a_{2} b_{2}, a_{2} b_{3}\right)$.
$-\operatorname{dim}\left(\left|\psi_{1}\right\rangle \otimes\left|\psi_{2}\right\rangle\right)=\operatorname{dim}\left(\left|\psi_{1}\right\rangle\right) \cdot \operatorname{dim}\left(\left|\psi_{2}\right\rangle\right)$
$-\operatorname{size}\left(\left|\psi_{1}\right\rangle \otimes\left|\psi_{2}\right\rangle\right)=\operatorname{size}\left(\left|\psi_{1}\right\rangle\right)+\operatorname{size}\left(\left|\psi_{2}\right\rangle\right)$
- size: number of qubits.
- Notation: $|0\rangle^{\otimes n}=|0\rangle \otimes \cdots \otimes|0\rangle, n$ times.


## Quantum mechanics in one slide



## Density matrix

- If a system is in state $\left|\psi_{1}\right\rangle$ with probability $p_{1}$, and in state $\left|\psi_{2}\right\rangle$ with probability $p_{2}$, then the system is in a mixed state.
- The mixed state is represented as a density matrix

$$
\rho=p_{1}\left|\psi_{1}\right\rangle\left\langle\psi_{1}\right|+p_{2}\left|\psi_{2}\right\rangle\left\langle\psi_{2}\right| .
$$

- In general, if a system is in state $\left|\psi_{i}\right\rangle$ with probability $p_{i}$, then the mixed state is

$$
\rho=\sum_{i} p_{i}\left|\psi_{i}\right\rangle\left\langle\psi_{i}\right|
$$

- For pure state $|\psi\rangle, \rho=|\psi\rangle\langle\psi|$.


## Density matrix

- Fact. A matrix $\rho$ is a density matrix of some mixed quantum state iff
$-\rho$ is positive semi-definite (psd)
$-\operatorname{Tr}(\rho)=1$.
- Recall:
- A matrix $M$ is psd if all its eigenvalues are nonnegative. Equivalently, if $\langle v| M|v\rangle \geq 0, \forall v$.
- The trace of a matrix $M$ is $\operatorname{Tr}(M)=\sum_{i} M_{i i}$.


## Postulates on mixed states

- Unitary operation $U: \rho \mapsto U \rho U^{\dagger}$
- For pure state $\rho=|\phi\rangle\langle\phi|$, it becomes $U \rho U^{\dagger}=$ $U|\phi\rangle\langle\phi| U^{\dagger}=\left|\phi^{\prime}\right\rangle\left\langle\phi^{\prime}\right|$ where $\left|\phi^{\prime}\right\rangle=U|\phi\rangle$.
- Orthogonal measurement $\left\{\left|\psi_{1}\right\rangle, \ldots\left|\psi_{d}\right\rangle\right\}$ : outcome $i$ occurs with probability $|\langle\psi| \rho| \psi\rangle\left.\right|^{2}$, and the system collapses to $\rho^{\prime}=\left|\psi_{i}\right\rangle\left\langle\psi_{i}\right|$.
- For pure state $\rho=|\phi\rangle\langle\phi|$, the probability is $\left|\left\langle\phi \mid \psi_{i}\right\rangle\right|^{2}$, and the collapsed state is $\left|\psi_{i}\right\rangle$.
- If we measure $\rho \in \mathbb{C}^{d \times d}$ in the computational basis $\{|1\rangle,|2\rangle, \ldots,|d\rangle\}$, then $\operatorname{Pr}\left[\right.$ outcome $i$ occurs] $=\rho_{i i}$, the $i$-th diagonal entry of $\rho$.
- Composition of $\rho_{1}$ and $\rho_{2}$ is just $\rho_{1} \otimes \rho_{2}$
- For pure state $\rho_{1}=\left|\phi_{1}\right\rangle\left\langle\phi_{1}\right|$ and $\rho_{2}=$ $\left|\phi_{2}\right\rangle\left\langle\phi_{2}\right|$, the joint state is

$$
\begin{aligned}
& \left|\phi_{1}\right\rangle\left\langle\phi_{1}\right| \otimes\left|\phi_{2}\right\rangle\left\langle\phi_{2}\right| \\
& \quad=\left(\left|\phi_{1}\right\rangle \otimes\left|\phi_{2}\right\rangle\right)\left(\left\langle\phi_{1}\right| \otimes\left\langle\phi_{2}\right|\right)
\end{aligned}
$$

- Recall tensor product of matrices:

$$
\left[\begin{array}{ll}
a_{11} & a_{12} \\
a_{21} & a_{22}
\end{array}\right] \otimes\left[\begin{array}{ll}
b_{11} & b_{12} \\
b_{21} & b_{22}
\end{array}\right]=\left[\begin{array}{llll}
a_{11} b_{11} & a_{11} b_{12} & a_{12} b_{11} & a_{12} b_{12} \\
a_{11} b_{21} & a_{11} b_{22} & a_{12} b_{21} & a_{12} b_{22} \\
a_{21} b_{11} & a_{21} b_{12} & a_{22} b_{11} & a_{22} b_{12} \\
a_{21} b_{21} & a_{21} b_{22} & a_{22} b_{21} & a_{22} b_{22}
\end{array}\right]
$$

- For operation, measurement and composition, these formulas for mixed states are all consistent to what we learned for pure states.
- So the formulas for mixed states extend those for pure states.


## entanglement

- Consider the following EPR pair state in a 2-qubit system: $\quad \frac{|00\rangle+|11\rangle}{\sqrt{2}}$
- It's in superposition of $|00\rangle$ and $|11\rangle$.
- There is no classical counterpart of this.
- Question: Is it really different than the classical correlation
$\left\{\begin{array}{ll}00 & \text { with prob. } 1 / 2 \\ 11 & \text { with prob. } 1 / 2\end{array} ?\right.$


## CHSH non-local game

- $x \in_{R}\{0,1\}, y \in_{R}\{0,1\}$
- Goal: A outputs $s$ and B outputs $t$ s.t.

$$
s \oplus t=x \cdot y
$$

- Value = largest $\operatorname{Pr}[s \oplus t=x \cdot y]$.
- Classical value: $3 / 4=0.75$.
- Even with shared randomness.
- Quantum value: $\frac{1}{2}+\frac{1}{2 \sqrt{2}} \approx 0.85$
$s \oplus t=x \cdot y ?$
- By sharing an EPR pair.


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- Intro to math model of quantum mechanics
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## Areas in quantum computing

- Quantum algorithms
- Quantum complexity
- Quantum cryptography
- Quantum error correction
- Quantum information theory
- Others: Quantum game theory / control /


## Area 1: Quantum Algorithms

## $\begin{array}{llllllll}1994 & 1996 & 1998 & 2000 & 2002 & 2004 & 2006 & 2008\end{array}$

## Shor: Factoring <br> \& Discrete Log

QFT
(Quantum Fourier Transform): exponential speedup; slower than expected.

- Factoring: Given an $n$-bit number, factor it (into product of two numbers).
- Classical (best known) : O( $\left.2^{n^{1 / 3}}\right)$
- Quantum*1: $O\left(n^{2}\right)$
*1: P. Shor. STOC"93, SIAM Journal on Computing, 1997.


## Area 1: Quantum Algorithms

## $199419961998 \quad 2000 \quad 2002 \quad 2004 \quad 2006 \quad 2008$


(Quantum Fourier Transform): exponential speedup; slower than expected.

- Implication of fast algorithm for Factoring
- Theoretical: Church-Turing thesis
- Practical: Breaking RSA-based cryptosystems


## Area 1: Quantum Algorithms

$\begin{array}{llllllll}1994 & 1996 & 1998 & 2000 & 2002 & 2004 & 2006 & 2008\end{array}$

## Shor: Factoring <br> \& Discrete Log

Hallgren: Pell's Equation
QFT
(Quantum Fourier Transform): exponential speedup; slower than expected.

- Pell's Equation: $x^{2}-n y^{2}=1$.
- Problem: Given $n$, find solutions $(x, y)$ to the above equation.
- Classical (best known):
$-\sim 2^{\sqrt{\log n}}$ (assuming the generalized Riemann hypothesis)
- ~ $n^{1 / 4}$ (no assumptions)
- Quantum*1: poly $(\log n)$.
*1: S. Hallgren. STOC'02. Journal of the ACM, 2007.


## Area 1: Quantum Algorithms

$\begin{array}{lllllllll}1994 & 1996 & 1998 & 2000 & 2002 & 2004 & 2006 & 2008\end{array}$


QFT
(Quantum Fourier Transform): exponential speedup; slower than expected.

- Hidden Subgroup Problem (HSP): Given a function $f$ on a group $G$, which has a hidden subgroup $H$, s.t. $f$ is
- constant on each coset $a H$,
- distinct on different cosets.

Task: find the hidden $H$.

- Factoring, Pell's Equation both reduce to it.
- Efficient quantum algorithms are known for Abelian groups.
- Main question: HSP for non-Abelian groups?


## Area 1: Quantum Algorithms

$\begin{array}{llllllll}1994 & 1996 & 1998 & 2000 & 2002 & 2004 & 2006 & 2008\end{array}$

| Shor: Factoring <br> \& Discrete Log | Hallgren: Pell's Equation | Kuperberg: HSP-Dihedral |
| :---: | :---: | :---: |

(Quantum Fourier Transform): exponential speedup; slower than expected.

- Two biggest cases:
- HSP for symmetric group $S_{n}$ : Graph Isomorphism reduces to it.
- HSP for dihedral group $D_{n}$ : Shortest Lattice Vector reduces to it.
- $\operatorname{HSP}\left(D_{n}\right)$ :
- Classical (best known): $2^{\log |G|}$
- Quantum ${ }^{* 1}: 2^{O(\sqrt{l o g}|G|)} 2^{O(\sqrt{\log |G|})}$
*1: G. Kuperberg. arXiv:quant-ph/0302112, 2003.


## Area 1: Quantum Algorithms

$$
\begin{array}{llllllll}
1994 & 1996 & 1998 & 2000 & 2002 & 2004 & 2006 & 2008
\end{array}
$$

| Shor: Factoring <br> \& Discrete Log | Hallgren: Pell's <br> Equation$\quad$Kuperberg: <br> HSP-Dihedral |
| :---: | :---: |

(Quantum Fourier Transform): exponential speedup; slower than expected.
Grover:
Search
(Quantum Search): polynomial speedup; most solved.

- Given $n$ bits $x_{1}, \ldots, x_{n}$, find an $i$ with $x_{i}=1$.
- Classical: $\Theta(n)$
- Quantum ${ }^{* 1}: ~ \Theta(\sqrt{n})$
*1: L. Grover. Physical Review Letters, 1997.


## Area 1: Quantum Algorithms

$\begin{array}{lllllllll}1994 & 1996 & 1998 & 2000 & 2002 & 2004 & 2006 & 2008\end{array}$

| Shor: Factoring <br> \& Discrete Log | Hallgren: Pell's <br> Equation |
| :---: | :---: |
| Kuperberg: <br> HSP-Dihedral |  |

(Quantum Fourier Transform): exponential speedup; slower than expected.

| Grover: |
| :--- |
| Search |

Many combinatorial /graph problems
(Quantum Search): polynomial speedup; most solved.

(Quantum Walk): poly and exp speedup; rapidly developed.
*1: D. Aharonov, A. Ambainis, J. Kempe, U. Vazirani. STOC'01

## Area 1: Quantum Algorithms



- Classical random walk on graphs: starting from some vertex, repeatedly go to a random neighbor
- Many algorithmic applications
- Quantum walk on graphs: even definition is non-trivial.
- For instance: classical random walk converges to a stationary distribution, but quantum walk doesn't since unitary is reversible.

(Quantum Walk): poly and exp speedup; rapidly developed.
*1: D. Aharonov, A. Ambainis, J. Kempe, U. Vazirani. STOC'01


## Area 1: Quantum Algorithms



- Element Distinctness: Given $n$ integers, decide whether they are the all distinct.
- Classical: $\Theta(n)$
- Quantum: $\Theta\left(n^{2 / 3}\right)$

(Quantum Walk): poly and exp speedup; rapidly developed.
*1: A. Ambainis, FOCS'04


## Area 1: Quantum Algorithms




Grover's search: OR function

general formula by \{AND-OR-NOT\}

- Classical: $\Theta(n)$
- Quantum: ~ $\Theta(\sqrt{n})$
- apply QW on the formula graph with weight carefully designed for inductions to work.

(Quantum Walk): poly and exp speedup; rapidly developed.
*1: A. Ambainis, A. Childs, B. Reichardt, R. Spalek, S. Zhang. FOCS'07


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## Power of quantum computing

- Question: How powerful is quantum computer?
- P: problems solvable in polynomial time
- One characterization of efficient computation
- BPP: problems solvable in probabilistic polynomial time w/ a small error tolerated
- Another characterization of efficient computation
- BQP: problems solvable in polynomial time by a quantum computer w/ a small error tolerated
- Yet another characterization of efficient computation, if you believe large-scale quantum mechanics.


## Classical upper bound of BQP

- Central in complexity theory: comparison of different modes of computation
- How to compare classical and quantum efficient computation?
- Quantum is more powerful: $\mathrm{BPP} \subseteq \mathrm{BQP}$
- An upper bound (of quantum by classical)
- [Thm*1] BQP $\subseteq$ PSPACE
- PSPACE: problems solvable in polynomial space.
- Believed to be much larger than NP.
*1: Bernstein, Vazirani. STOC'93, SIAM J. on Computing, 1997


## Where does BQP sit in?

- BQP contains BPP and $P$.
- But it probably doesn't contain all NP.
- Yet it's possible to be outside PH.
- It's position may be weird.


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## Meyer's game: classical



- Actions $A_{i}, B_{i}$ : to flip or not to flip
- Alice's Goal: 0. Bob’s Goal: 1.
- A Nash equilibrium: $A_{i}, B_{i}$ flip with half prob.
- Then each wins with half prob.


## Meyer's game: quantum



- Bob remains classical: $B_{1}$ is either $X=\left[\begin{array}{cc}0 & 1 \\ 1 & 0\end{array}\right]$ (Swap $|0\rangle$ and $|1\rangle$ ) or identity (doing nothing).
- Alice is quantum: $A_{i}$ can be any 1 -qubit operation.
- Alice's Goal: |0才. Bob's Goal: |1>.
- Now Alice can win for sure by applying a Hadamard gate. $A_{1}:|0\rangle \rightarrow|+\rangle . A_{2}:|+\rangle \rightarrow|0\rangle$.


## Meyer's game: fairness issue



- Despite the quantum advantage, there is clear a fairness issue.
- Alice has two actions.
- And the actions are in a fixed order of $A_{1} \rightarrow B_{1} \rightarrow A_{2}$.
- Question: Can quantum advantage still exist in a more fair setting?
- For fairness: each player makes just one action, simultaneously.
- This is nothing but strategic games!


## Quantization*1 of strategic game: Penny Matching

 potential action classical outcome utility

- $|\varphi\rangle$ is an equilibrium if both players are classical,
- Each wins with half prob.
- If Alice turns to quantum: $A=H$ turns $|\varphi\rangle$ into $\frac{|0\rangle|0\rangle+|1\rangle|1\rangle}{\sqrt{2}}$. Then she wins for sure!
- Message: quantum player has a huge advantage when playing against a classical player.
*1. Zu, Wang, Chang, Wei, Zhang, Duan, NJP, 2012.

Quantization of strategic game: Penny Matching potential action classical outcome utility


- State is symmetric, so it doesn't matter who takes which qubit.
- We can also let the classical player Bob to choose the target goal.
- If Bob wants $a=b$, then Alice applies $X H$.


## Quantum advantage in strategic games

potential action classical outcome utility


No! $\rho=\frac{1}{4}\binom{+\mid+\langle+| \otimes|0\rangle\langle 0|+|0\rangle\langle 0| \otimes|+\rangle\langle+|}{+|-\rangle\langle-| \otimes|1\rangle\langle 1|+|1\rangle\langle 1| \otimes|-\rangle\langle-|} \quad \begin{aligned} & \text { No entanglement. } \\ & \text { But has discord. }\end{aligned}$

- $\rho$ is an equilibrium if both players are classical, each winning with prob. $=1 / 2$
- If Alice uses quantum, $A=H$ increases her winning prob. to $3 / 4$.
- Question*2: Is discord necessary?
- Yes, if each player's part (of the shared state) is a qubit,
- No, if each player's part (of the shared state) has dimension 3 or more.
*2. Wei, Zhang, TAMC, 2014.


## Games between quantum players

- After these examples, Bob realizes that he should use quantum computers as well.
- Question: Any advantage when both players are quantum?
- Previous correspondence results imply a negative answer for complete information games.
- But quantum advantage exists for Bayesian games!


## Quantum Bayesian games



- Each player $i$ has a private input/type $x_{i}$.
- $x_{i}$ is known to Player $i$ only.
- The joint input is drawn from some distribution $P$.
- Each player $i$ can potentially apply some operation $\Phi_{i}$.
- A measurement in the computational basis gives output $\left|y_{i}\right\rangle$ for Player $i$, who receives utility $u_{i}\left(x_{1}, x_{2}, y_{1}, y_{2}\right)$.


## Quantum Bayesian games

potential action classical outcome
utility


- Classical state $|\varphi\rangle=\left(r_{1}, r_{2}\right) \leftarrow$ distribution $Q$.
- Classical strategy $\Phi_{i}=c_{i}\left(x_{i}, r_{i}\right)$.
- Classical payoff

$$
\mathbf{E}\left[u_{i}\right]=\mathbf{E}_{x \leftarrow P, r \leftarrow Q}\left[u_{i}\left(x, c_{1}\left(x_{1}, r_{1}\right), c_{2}\left(x_{2}, r_{2}\right)\right]\right.
$$

- $\left(Q, c_{1}, c_{2}\right)$ is equilibrium if no player can gain a higher payoff by changing her strategy unilaterally.


## Quantum Bayesian games

potential action classical outcome
utility


- Quantum strategy $\Phi_{1}=\left\{E_{x_{1}}^{y_{1}}: E_{x_{1}}^{y_{1}} \succcurlyeq 0, \sum_{y_{1}} E_{x_{1}}^{y_{1}}=I\right\}, \Phi_{2}=$ $\left\{F_{x_{2}}^{y_{2}}: F_{x_{2}}^{y_{2}} \succcurlyeq 0, \sum_{y_{2}} F_{x_{2}}^{y_{2}}=I\right\}$.
- Quantum payoff

$$
\mathbf{E}\left[u_{i}\right]=\mathbf{E}_{x \leftarrow P}\left[\langle\varphi| E_{x_{1}}^{y_{1}} \otimes F_{x_{2}}^{y_{2}}|\varphi\rangle \cdot u_{i}(x, y)\right]
$$

- $\left(|\varphi\rangle, \Phi_{1}, \Phi_{2}\right)$ is equilibrium if no player can gain a higher payoff by changing her strategy unilaterally.


## Quantum Bayesian games


*1. Pappa, Kumar, Lawson, Santha, Zhang, Diamanti, Kerenidis, PRL, 2015.

## Quantum Bayesian games

potential actionclassical outcome
utility

|  | $y_{B}=0$ | $y_{B}=1$ |
| :---: | :---: | :---: |
| $y_{A}=0$ | $(0,0)$ | $(3 / 4,3 / 4)$ |
| $y_{A}=1$ | $(3 / 4,3 / 4)$ | $(0,0)$ |

Table II: $x_{A} \wedge x_{B}=1$
*1. Pappa, Kumar, Lawson, Santha, Zhang, Diamanti, Kerenidis, PRL, 2015.

## Viewed as non-locality

- Traditional quantum non-local games exhibit quantum advantages when the two players have the common goal.
- CHSH, GHZ, Magic Square Game, Hidden Matching Game, Brunner-Linden game.
- Now the two players have conflicting interests.
- Quantum advantages still exist.
- Message: If both players play quantum strategies in an equilibrium, they can also have advantage over both being classical.


## Summary

- Quantum algorithms: offer huge speedup for certain computational problems.
- Quantum entanglement:
- A distinctive feature of quantum mechanics.
- Proof that our world is quantum mechanical.
- Quantum games: quantum players can have big advantages.

